

The two-dimensional form of Kirchhoff's formula (Papoulis, p. 302)

$$1. \quad v(x_0, z_0) = \frac{1}{4} \int_{-\infty}^{\infty} \left[ u(x, 0) \frac{\partial H_0(k\rho)}{\partial n} - H_0(k\rho) \frac{\partial u(x, 0)}{\partial n} \right] dx$$

with

$$2. \quad \rho = \sqrt{(x-x_0)^2 + (z-z_0)^2}$$

and  $H_0(k\rho)$  the Hankel function of the first kind order zero.

According to Papoulis for  $k\rho \gg 1$

$$3. \quad H_0(k\rho) \sim \sqrt{\frac{2j}{\pi k\rho}} e^{jk\rho}$$

with  $j$  the imaginary unit.

Furthermore,  $k = 2\pi/\lambda$  with  $k$  the wavenumber and  $\lambda$  the wavelength.

$$4. \quad v(x_0, z_0) = \sqrt{\frac{j}{8\pi k}} \int_{-\infty}^{\infty} \left[ u(x, 0) \frac{\partial}{\partial n} \left( \frac{e^{jk\rho}}{\sqrt{\rho}} \right) - \frac{e^{jk\rho}}{\sqrt{\rho}} \frac{\partial u(x, 0)}{\partial n} \right] dx$$

We assume that

$$5. \quad u(x, z) = u(x) e^{jkz} \quad \text{and} \quad \frac{\partial}{\partial n} = \frac{\partial}{\partial z}$$

In view of  $\frac{\partial F(\rho(z))}{\partial z} = \frac{dF}{d\rho} \frac{d\rho}{dz}$  we find

$$6. \quad \frac{\partial}{\partial z} \left( \frac{e^{jk\rho}}{\sqrt{\rho}} \right) = \frac{e^{jk\rho}}{\sqrt{\rho}} \left( \frac{jk(z-z_0)}{\rho} - \frac{z-z_0}{z\rho^2} \right) \sim \frac{e^{jk\rho}}{\sqrt{\rho}} \frac{jk(z-z_0)}{\rho}$$

Because  $k(z-z_0)/\rho \gg (z-z_0)/(z\rho^2)$  for  $k\rho \gg 1$ .

$$7. \quad \frac{\partial u(x, 0)}{\partial n} = jk u(x)$$

$$8. \quad v(x_0, z_0) = \sqrt{\frac{j}{8\pi k}} \int_{-\infty}^{\infty} u(x) \left[ -\frac{jkz_0}{\rho} - jk \right] \frac{e^{jk\rho}}{\sqrt{\rho}} dx$$

We can write this formula as

$$9. \quad v(x_0, z_0) = \int_{-\infty}^{\infty} u(x) g(\rho(x)) e^{jk\rho(x)} dx$$

with

$$10. \quad g(\rho(x)) = \frac{1}{2\sqrt{j\lambda\rho}} \left( 1 + \frac{z_0}{\rho} \right)$$

We assume that

$$11. \quad u(x) = \int_{-\infty}^{\infty} u(f) e^{j2\pi fx} df$$

$$12. \quad v(x_0, z_0) = \int_{-\infty}^{\infty} u(f) \left[ \int_{-\infty}^{\infty} g(\rho(x)) e^{jk\rho(x) + j2\pi fx} dx \right] df$$

We define

$$13. \quad I(f) = \int_{-\infty}^{\infty} g(\rho(x)) e^{jk\rho(x) + j2\pi fx} dx$$

$$14. \quad \Phi(x) = 2\pi fx + k\rho(x)$$

With  $x_{sf}$  the point of stationary phase we can write

$$15. \quad \Phi(x) = \Phi(x_{sf}) + (x-x_{sf})\Phi'(x_{sf}) + \frac{(x-x_{sf})^2}{2!}\Phi''(x_{sf}) + \dots$$

For the point  $x_{sf}$  we know that,

$$16. \quad \Phi'(x_{sf}) = 0$$

so we find, with  $z=0$ , the following formulae.

$$17. \quad x_{sf} = x_0 - \frac{\lambda f z_0}{\sqrt{1-\lambda^2 f^2}}$$

$$18. \quad \rho(x_{sf}) = \frac{z_0}{\sqrt{1-\lambda^2 f^2}}$$

$$19. \quad \Phi(x_{sf}) = 2\pi f x + k z_0 \sqrt{1-\lambda^2 f^2}$$

$$20. \quad \Phi''(x_{sf}) = \frac{k}{z_0} (1-\lambda^2 f^2)^{3/2}$$

This leads to the following formula

$$21. \quad I(f) = g(\rho(x_{sf})) e^{i\Phi(x_{sf})} \int_{-\infty}^{\infty} e^{i\frac{k}{2z_0}(1-\lambda^2 f^2)^{3/2}(x-x_{sf})^2} dx$$

For the Gaussian integral it is known that

$$22a. \quad \int_{-\infty}^{\infty} e^{i\alpha y^2} dy = \sqrt{\pi} e^{i\pi/4} \quad \text{and consequently } 22b. \quad \int_{-\infty}^{\infty} e^{i\alpha y^2} dy = \sqrt{\frac{\pi}{\alpha}}$$

which leads to

$$23. \quad I(f) = \frac{1}{2} \left(1 + \frac{1}{\sqrt{1-\lambda^2 f^2}}\right) e^{i k z_0 \sqrt{1-\lambda^2 f^2}} e^{i 2\pi f x}$$

so we conclude that

$$24. \quad v(x_0, z_0) = \int_{-\infty}^{\infty} u(f) \left[ \frac{1}{2} \left(1 + \frac{1}{\sqrt{1-\lambda^2 f^2}}\right) e^{i k z_0 \sqrt{1-\lambda^2 f^2}} \right] e^{i 2\pi f x} dx$$

We define the optical transfer function  $H_{FK}(f)$

$$25. \quad v(x_0, z_0) = \int_{-\infty}^{\infty} u(f) H_{FK}(f) e^{i 2\pi f x} df$$

(F=Fresnel and K=Kirchhoff).

$$26. \quad H_{FK}(f) = \frac{1}{2} \left(1 + \frac{1}{\sqrt{1-\lambda^2 f^2}}\right) e^{i k z_0 \sqrt{1-\lambda^2 f^2}}$$

It is well-known that (Goodman, p. 60).

$$27. \quad H_{RS1}(f) = e^{i k z_0 \sqrt{1-\lambda^2 f^2}}$$

(R=Rayleigh and S=Sommerfeld).

We conclude that

$$28. \quad H_{RS2}(f) = \frac{1}{\sqrt{1-\lambda^2 f^2}} e^{i k z_0 \sqrt{1-\lambda^2 f^2}}$$

because the Fresnel-Kirchhoff solution is the average of the two Rayleigh-Sommerfeld solutions (Goodman, p. 50).

References.

J.W. Goodman, Introduction to Fourier Optics, 1996.

A. Papoulis, Systems and Transforms with Applications in Optics, 1968.