

$e^{n\alpha}$  in the expression representing  $\psi$ . The values so obtained for the coefficients in (83) above are given by

$$A_1 = \frac{c \sin 2\beta_0 - 2\beta_0 \cos 2\beta_0}{\pi (2\beta_0 - \sin 2\beta_0)} \dots \dots \dots (88)$$

$$(A_1 + A_2) = \frac{2c \sin 2\beta_0 - 2\beta_0 \cos 2\beta_0}{\pi (4\beta_0 - \sin 4\beta_0)} \dots \dots \dots (89)$$

$$(A_2 + A_3) = (A_1 + A_2) \frac{c \frac{1}{2} \sin 4\beta_0 - \sin 2\beta_0}{\pi \frac{1}{2} (6\beta_0 - \sin 6\beta_0)} \dots \dots \dots (90)$$

$$A_n + A_{n+1} = (A_{n-1} + A_n) \frac{c \frac{1}{n} \sin 2n\beta_0 - \sin 2\beta_0}{\pi \frac{1}{n+1} (2(n+1)\beta_0 - \sin 2(n+1)\beta_0)} \dots \dots (91)$$

Of the constants  $A'_0, A'_1$ , etc.  $A'_0 = -3\eta_0^2$ ; and the remaining coefficients are to be determined by equating to zero the expressions multiplying each  $e^{n\alpha}$  in (87) above. Thus the solutions representing flow into a parallel channel within the boundaries specified by  $\eta = \eta_0$  have been found; and the fluid movement represented is unique in accordance with the well-known theorem of Helmholtz (6). The study of these solutions in detail must be deferred to a later opportunity.

References.

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XXXII. Coefficients for Numerical Integration with Central Differences.

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THE following table lists the coefficients in the formula for numerical integration employing central differences through the fifty-fifth order. Previous calculations have gone only as far as coefficients of the seventh central difference. The coefficients  $M_{2s}$  occur in the following formula :

$$\frac{1}{h} \int_a^{a+nh} f(x) dx = [\frac{1}{2}f(a) + f(a+h) + \dots + f(a+(n-1)h) + \frac{1}{2}f(a+nh)] + \sum_{s=1}^{n-1} M_{2s} [\mu \delta^{2s-1} f(a+nh) - \mu \delta^{2s-1} f(a)] + nh^{2m} M_{2m} f^{(2m)}(\xi). \quad (1)$$

\* Communicated by the Author.

† Mathematical Tables Project, National Bureau of Standards, U.S.A.

The coefficient (2s)th Bernoulli function (1) is known Gaussian sum coefficients, given

where  $D^{2r+1}O(x)$  polynomial

for  $x=0$ .

which is identical

For large  $s$ ,

(At  $s=10$ , this

$$M_{2s} = \frac{1^2}{3! 2}$$

The first exact values checked both

2195  
2196

The coefficient  $M_{2s}$  is identical with  $B_{2s}^{(2s)}(s)/(2s)!$  where  $B_{2s}^{(2s)}(s)$  is the  $(2s)$ th Bernoulli polynomial of order  $2s$  for argument equal to  $s$ . Equation (1) is known as the Gauss-Encke formula and also as the second Gaussian summation-formula. The following properties of these coefficients, given in (2)-(7), are all demonstrated in Steffensen :

$$M_{2s+1} = 0, \dots \dots \dots (2)$$

$$M_{2s} = \frac{2}{(2s)!} \sum_{r=0}^{r=s} \frac{D^{2r+1}O^{(2s+1)}}{2^{2r+1}(2r+1)(2r+1)!}, \dots \dots (3)$$

where  $D^{2r+1}O^{(2s+1)}$  is the  $(2r+1)$ th derivative of the central factorial polynomial

$$x^{(2s+1)} \equiv x(x^2 - 1/4)(x^2 - 9/4) \dots \left[ x^2 - \frac{(2s-1)^2}{4} \right]$$

for  $x=0$ .

$$M_{2s} = \frac{2}{(2s)!} \int_0^1 \frac{\Gamma(x+s+\frac{1}{2})}{\Gamma(x-s+\frac{1}{2})} dx, \dots \dots (4)$$

which is identical with Milne-Thomson's

$$M_{2s} = \frac{1}{(2s)!} \int_0^1 (x+s-1)(x+s-2) \dots (x-s) dx, \dots \dots (4a)$$

$$|M_{2s}| \leq \binom{2s}{s} 2^{-4s}. \dots \dots (5)$$

For large  $s$ ,

$$M_{2s} \sim \frac{(-1)^s}{2^{2s-1} s^{3/2} \pi^{3/2}} \dots \dots (6)$$

(At  $s=10$ , this formula is accurate to about 1 per cent.)

$$M_{2s} = \frac{1^2}{3! 2^2} M_{2s-2} - \frac{(1.3)^2}{5! 2^4} M_{2s-4} + \frac{(1.3.5)^2}{7! 2^6} M_{2s-6} - \dots$$

$$+ \frac{(-1)^s [1.3 \dots (2s-3)]^2}{(2s-1)! 2^{2s-2}} M_2 + \frac{(-1)^s [1.3 \dots (2s-1)]^2}{(2s-1)! (2s+1) 2^{2s}}. \quad (7)$$

The first twenty coefficients given below were checked, using their exact values, in the cumulative recursion formula (7). All values were checked both by differencing the ratios  $M_{2s}/M_{2s+2}$  and by actual integrations.

Table of Coefficients.

$M_2$	$-\frac{1}{12}$
$M_4$	$+\frac{11}{720}$
$M_6$	$-\frac{191}{60480}$

numerators = #2195  
denominators = #2196

$M_8$	$+ \frac{2497}{3628800}$
$M_{10}$	$- \frac{14797}{95800320}$
$M_{12}$	$+ \frac{924 \quad 27157}{261 \quad 53487 \quad 36000}$
$M_{14}$	$- \frac{367 \quad 40617}{448 \quad 34549 \quad 76000}$
$M_{16}$	$+ \frac{6 \quad 14309 \quad 43169}{32 \quad 01186 \quad 85286 \quad 40000}$
$M_{18}$	$- \frac{2313 \quad 39458 \quad 92303}{51090 \quad 94217 \quad 17094 \quad 40000}$
$M_{20}$	$+ \frac{1639 \quad 96886 \quad 81447}{152579 \quad 28431 \quad 37024 \quad 00000}$
$M_{22}$	$- 00000 \quad 00256 \quad 38286 \quad 986$
$M_{24}$	$- 00000 \quad 00061 \quad 40295 \quad 342$
$M_{26}$	$- 00000 \quad 00014 \quad 75583 \quad 052$
$M_{28}$	$- 00000 \quad 00003 \quad 55628 \quad 077$
$M_{30}$	$- 00000 \quad 00000 \quad 85924 \quad 015$
$M_{32}$	$- 00000 \quad 00000 \quad 20805 \quad 605$
$M_{34}$	$- 00000 \quad 00000 \quad 05047 \quad 538$
$M_{36}$	$- 00000 \quad 00000 \quad 01226 \quad 642$
$M_{38}$	$- 00000 \quad 00000 \quad 00298 \quad 549$
$M_{40}$	$- 00000 \quad 00000 \quad 00072 \quad 762$
$M_{42}$	$- 00000 \quad 00000 \quad 00017 \quad 755$
$M_{44}$	$- 00000 \quad 00000 \quad 00004 \quad 338$
$M_{46}$	$- 00000 \quad 00000 \quad 00001 \quad 061$
$M_{48}$	$- 00000 \quad 00000 \quad 00000 \quad 260$
$M_{50}$	$- 00000 \quad 00000 \quad 00000 \quad 064$
$M_{52}$	$- 00000 \quad 00000 \quad 00000 \quad 016$
$M_{54}$	$- 00000 \quad 00000 \quad 00000 \quad 004$
$M_{56}$	$- 00000 \quad 00000 \quad 00000 \quad 001$

— end

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XXXIII. *Equilib.*

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- (i.) "Westergaard proportional to the i
  - (ii.) A note by in infinite extent restin
  - (iii.) A paper gi corresponding to (ii.)

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